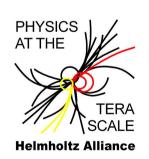
- HELAC-NLO -Developments and Applications





Outline

- ☐ General motivation for NLO QCD calculations
- **HELAC-NLO** in a nutshell
- \square Applications: 2 -> 4 processes:
 - ttbb
 - ttjj
 - WWbb
- ☐ Summary & Outlook

HELAC-NLO Group:

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Introduction

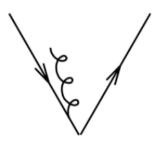
- 8-10 partons in the final state @ LO, well separated to avoid divergences
- ☐ On the market automatic parton level tools which are completely self contained
- ☐ Provide amplitudes and integrators on their own
- Standard Model and beyond tools @ tree level (just few examples)
 - > ALPGEN, AMEGIC++/SHERPA, COMIX/SHERPA, HELAC-PHEGAS, MADGRAPH/MADEVENT, O'MEGA/WHIZARD, ...
- General purpose Monte Carlo programs (parton shower, hadronization, multiple interactions, hadrons decays, etc.)
 - ➤ HERWIG, HERWIG++, PYTHIA 6.4, PYTHIA 8.1, SHERPA, ...
- High sensitivity to unphysical input scales, to improve accuracy of prediction higher order calculations are needed

Motivation for NLO

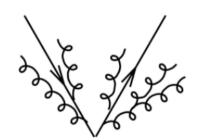
- Stabilizing the scale in the QCD input parameters most notably the strong coupling constant and PDFs
- Normalization and shape of distributions first known at NLO
- Many scale processes: V+ jets, VV + jets, ttH, tt + jets, njets ...
- Sometimes dynamical scales seem to work better for some observables
- How do we know which scale to choose?
- ☐ Improved description of jets



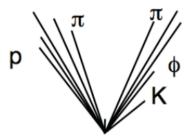




NLO



Parton Shower



Hadron Level

Structure of NLO Calculations

$$\sigma^{NLO} = \int_{m} d\sigma^{B} + \int_{m+1} d\sigma^{R} - \int_{m+1} d\sigma^{A} + \int_{m+1} d\sigma^{A} + \int_{m} d\sigma^{V}$$

$$\hookrightarrow \int d\sigma^{B} + \int_{m+1} \left[d\sigma^{R} - d\sigma^{D} \right] + \int_{m} \left[d\sigma^{V} + d\sigma^{I} + d\sigma^{KP} \right]$$

- ☐ Our strategy in a few words
 - > make it fully numeric
 - > make it fully automatic
 - > ,,montecarlize" everything for speed

☐ Decompose the amplitude into a basis of scalar integrals

$$\mathcal{A} = \sum d_{i_1 i_2 i_3 i_4} + \sum c_{i_1 i_2 i_3} + \sum b_{i_1 i_2} + \sum a_{i_1} + \sum a_{i_1} + \sum a_{i_2} + \sum a_{i_3} + \sum a_{i_4} + \sum a_{i_5} + \sum a_{i_5}$$

$$\mathcal{A} = \sum_{I \subset \{0,1,\cdots,m-1\}} \int \frac{\mu^{(4-d)d^d q}}{(2\pi)^d} \frac{\bar{N}_I(\bar{q})}{\prod_{i \in I} \bar{D}_i(\bar{q})}$$

- ☐ Three main building blocks are needed
 - Evaluation of numerator function N(q)
 - > Determination of coefficients via reduction method
 - Evaluation of scalar functions via ONELOOP

- □ Reduction at integrand level **OPP** method implemented in **CUTTOOLS**
- ☐ Computing numerator functions for specific values of loop momentum that are solutions of equations

$$D_i(q) = 0$$
 for $i = 0, ..., M-1$

☐ It is customary to refer to these equations as quadruple (M = 4), triple (M = 3), double (M = 2) and single (M = 1) cuts

$$N(q) = \sum_{i_0 < i_1 < i_2 < i_3}^{m-1} \left[d(i_0 i_1 i_2 i_3) + \tilde{d}(q; i_0 i_1 i_2 i_3) \right] \prod_{i \neq i_0, i_1, i_2, i_3}^{m-1} D_i$$

$$+ \sum_{i_0 < i_1 < i_2}^{m-1} \left[c(i_0 i_1 i_2) + \tilde{c}(q; i_0 i_1 i_2) \right] \prod_{i \neq i_0, i_1, i_2}^{m-1} D_i$$

$$+ \sum_{i_0 < i_1}^{m-1} \left[b(i_0 i_1) + \tilde{b}(q; i_0 i_1) \right] \prod_{i \neq i_0, i_1}^{m-1} D_i$$

$$+ \sum_{i_0}^{m-1} \left[a(i_0) + \tilde{a}(q; i_0) \right] \prod_{i \neq i_0}^{m-1} D_i$$

$$+ \tilde{P}(q) \prod_{i=1}^{m-1} D_i.$$

- \square Compute the rational terms $R = R_1 + R_2$
- \square R₁ originates from ϵ dependence of denominators

$$D_i \to \bar{D}_i - \tilde{q}^2$$

Computed within the framework of OPP reduction

Ossola,, Papadopoulos, Píttau '07, '08

 \square R₂ originates from ϵ dependence of numerators

$$\bar{q} = q + \tilde{q}$$
 $\bar{\gamma}_{\mu} = \gamma_{\mu} + \tilde{\gamma}_{\mu}$ $\bar{g}^{\mu\nu} = g^{\mu\nu} + \tilde{g}^{\mu\nu}$

> Computed with effective tree-level Feynman rules

Draggiotis, Garzelli, Papadopoulos, Pittau '09 Garzelli, Malamos, Pittau '09

☐ In case of ttbb final state the integrand has the form

$$\mathcal{A}(q) = \sum \frac{N_{i}^{(6)}(q)}{\bar{D}_{i_{0}}\bar{D}_{i_{1}}\cdots\bar{D}_{i_{5}}} + \underbrace{\frac{N_{i}^{(5)}(q)}{\bar{D}_{i_{0}}\bar{D}_{i_{1}}\cdots\bar{D}_{i_{4}}}}_{-\bar{D}_{i_{0}}\bar{D}_{i_{1}}\cdots\bar{D}_{i_{3}}} + \underbrace{\frac{N_{i}^{(4)}(q)}{\bar{D}_{i_{0}}\bar{D}_{i_{1}}\cdots\bar{D}_{i_{3}}}}_{-\bar{D}_{i_{0}}\bar{D}_{i_{1}}\bar{D}_{i_{2}}} + \cdots$$

- \blacksquare **HELAC-1LOOP** evaluates numerically the numerators $N_i^6(q)$, $N_i^5(q)$, ...
- with the values of the loop momentum q provided by **CUTTOOLS**
 - ➤ Generates all partitions of 6, 5, 4, ... blobs attached to the loop and checks all possible flavors (colors) that can run inside
 - ➤ Hard cuts the loop (q is fixed) to get n+2 tree-like process

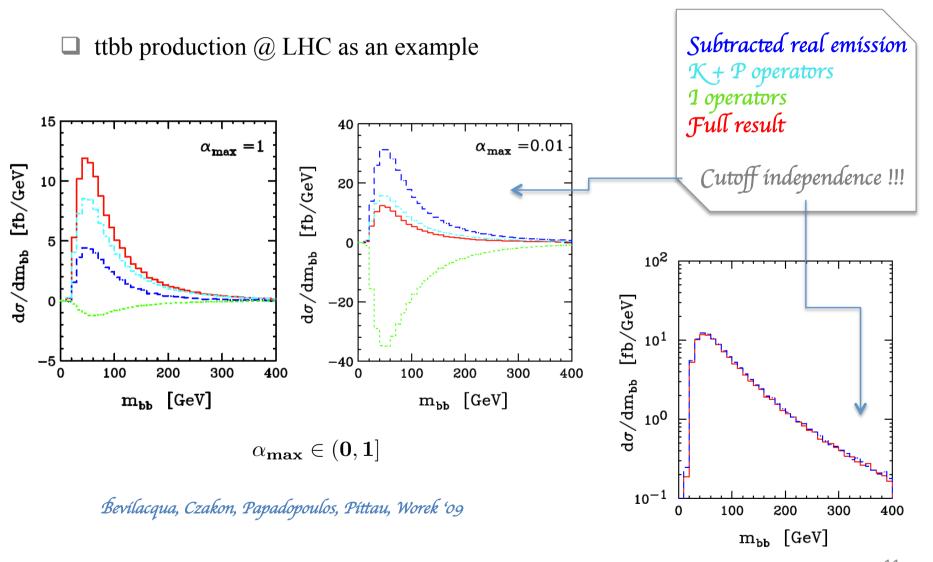
Van Hameren, Papadopoulos, Píttau '09 Bevílacqua, Czakon, Papadopoulos, Píttau, Worek '09

Real Emission

- ☐ HELAC-DIPOLES complete, public, automatic Catani-Seymour dipoles
- ☐ Phase space integration of subtracted real radiation and integrated dipoles in both massless and massive cases
- Extended for arbitrary polarizations
 - ➤ Monte Carlo over polarization states of external particles
 - Monte Carlo over color
- ☐ Phase space restriction on the dipoles phase space
 - Less dipole subtraction terms per event
 - > Increased numerical stability
 - Reduced missed binning problem
 - Large cancellations between subtracted real radiation and integrated dipoles

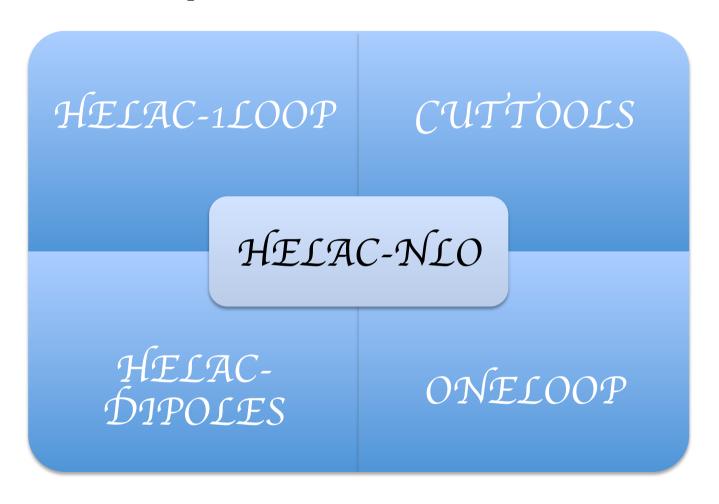
Czakon, Papadopoulos, Worek '09

Real Emission



Evaluation of loop numerators N(q) and R_2

Reduction of tensor integrals, OPP coefficients and R_1



Catani-Seymour dipole subtraction for massless and massive cases

Evaluation of scalar integrals

Applications

pp -> ttbb & pp -> ttjj

☐ Integrated cross sections and scale dependence Permille level agreement!

Process	$\sigma_{[23, 24]}^{LO}$ [fb]	σ^{LO} [fb]	$\sigma_{[23, 24]}^{\rm NLO} [{\rm fb}]$	$\sigma_{\alpha_{\text{max}}=1}^{\text{NLO}}$ [fb]	$\sigma_{\alpha_{\rm max}=0.01}^{\rm NLO}$ [fb]
$q\bar{q} \to t\bar{t}b\bar{b}$	85.522(26)	85.489(46)	87.698(56)	87.545(91)	87.581(134)
$pp \to t\bar{t}b\bar{b}$	1488.8(1.2)	1489.2(0.9)	2638(6)	2642(3)	2636(3)

$\xi \cdot m_t$	$1/8 \cdot m_t$	$1/2 \cdot m_t$	$1 \cdot m_t$	$2 \cdot m_t$	$8 \cdot m_t$
$\sigma^{\rm LO}$ [fb]	8885(36)	2526(10)	1489.2(0.9)	923.4(3.8)	388.8(1.4)
$\sigma^{\rm NLO}$ [fb]	4213(65)	3498(11)	2636(3)	1933.0(3.8)	1044.7(1.7)

$$\sigma_{
m LO} = (1489.2 \pm 0.9) {
m \ fb}$$

$$\sigma_{\mathbf{NLO}} = (\mathbf{2636} \pm \mathbf{3}) \ \mathbf{fb}$$

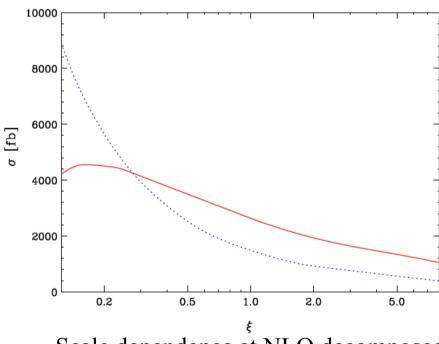
Bevilacqua, Czakon, Papadopoulos, Pittau, Worek '09 Bredenstein, Denner, Dittmaier, Pozzorini '08, '09 Scale dependence reduced:

70% (a) LO down to **33% (a) NLO**

K factor of K = 1.77

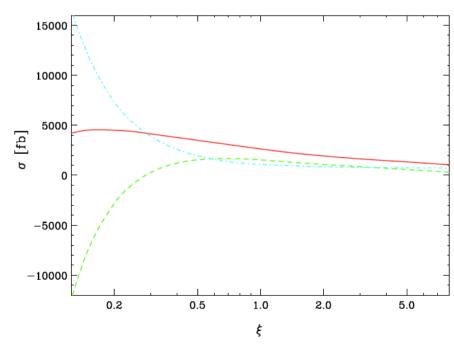
for quarks initial states only K = 1.03With jet veto of 50 GeV K = 1.20

☐ Scale dependence graphically



Scale dependence at NLO decomposed into contribution of *Virtual Corrections* & Real Radiation

Varying scale up or down by a factor two changes cross section by 70% (a) LO and by 33% (a) NLO

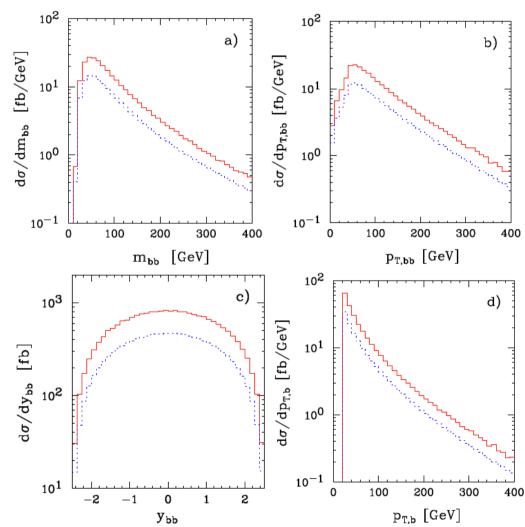


Bevilacqua, Czakon, Papadopoulos, Pittau, Worek '09

- Differential cross sections
- □ b-jet pair kinematics
 - > Invariant mass
 - Transverse momentum
 - Rapidity distribution
- □ single b-jet kinematics
 - > Transverse momentum

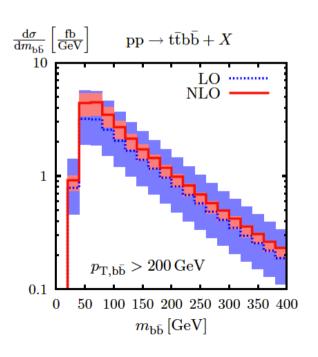
LO & NLO

■ Relatively small variation compared to the size but shape change important



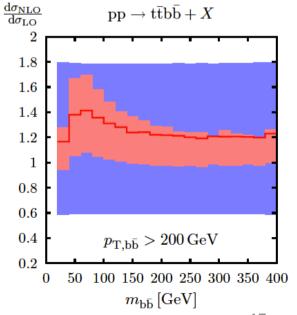
- Broad study
- Cross section in fb
- Dynamic scale
- \rightarrow m_{66} distribution
- > K-factor

$pp \rightarrow t\bar{t}b\bar{b} + X$ σ [fb] 3500 ro NLO 3000 $\mu_{\mathrm{R}}^2 = m_{\mathrm{t}} \sqrt{p_{\mathrm{T,b}} p_{\mathrm{T,\bar{b}}}} \xi^2$ 2500 $\mu_{\mathrm{F}}^2 = m_{\mathrm{t}} \sqrt{p_{\mathrm{T,b}} p_{\mathrm{T,\bar{b}}}} \, \xi^2$ 2000 $p_{\mathrm{T,b\bar{b}}} > \dot{200}\,\mathrm{GeV}$ 15001000 500 $0.125 \ 0.25$ 0.5ξ



Bredenstein, Denner, Dittmaier, Pozzorini '10

Setup	$m_{ m bar{b},cut}$	$p_{\mathrm{T,bar{b},cut}}$	$p_{ m jet,veto}$	$p_{\mathrm{T,b,cut}}$	$y_{ m b,cut}$	$\sigma_{ m LO}$	$\sigma_{ m NLO}$	K
Ι	100	-	-	20	2.5	$786.3(2)_{-41\%}^{+78\%}$	$978(3)_{-21\%}^{+13\%}$	1.24
II	-	200	-	20	2.5	$451.8(2)_{-41\%}^{+79\%}$	$592(4)_{-22\%}^{+13\%}$	1.31
III	100	-	100	20	2.5	$786.1(6)_{-41\%}^{+78\%}$	$700(3)_{-19\%}^{+0.4\%}$	0.89
IV	100	-	-	50	2.5	$419.4(1)_{-40\%}^{+77\%}$	$526(2)_{-21\%}^{+13\%}$	1.25

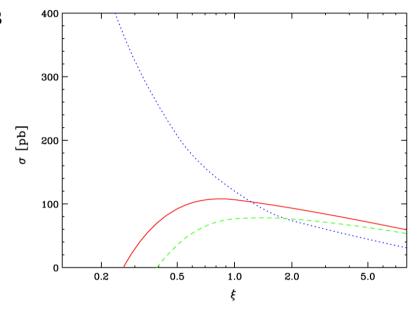


pp -> ttjj @ LHC

☐ Scale dependence & integrated cross sections

Bevilacqua, Czakon, Papadopoulos, Worek '10

Process	$\sigma^{ m LO}$ [pb]	Contribution
$pp \to t\bar{t}jj$	120.17(8)	100%
$qg \rightarrow t\bar{t}qg$	56.59(5)	47.1%
$gg \rightarrow t\bar{t}gg$	52.70(6)	43.8%
$qq' \to t\bar{t}qq', \ q\bar{q} \to t\bar{t}q'\bar{q}'$	7.475(8)	6.2%
$gg \to t\bar{t}q\bar{q}$	1.981(3)	1.6%
$q\bar{q} \to t\bar{t}gg$	1.429(1)	1.2%



 $\sigma_{\rm LO} = (120.17 \pm 0.08)~\rm pb$

$$\sigma_{
m NLO} = (106.94 \pm 0.17) \
m pb$$

$$\sigma_{\mathbf{NLO}}^{\mathbf{veto}} = (\mathbf{76.58} \pm \mathbf{0.17}) \ \mathbf{pb}$$

Scale dependence reduced:

72% @ LO down to 13% @ NLO 54% @ NLO with jet veto of 50 GeV

K factor of K = 0.89 (K = 0.64) Negative shift of 11% (36%)

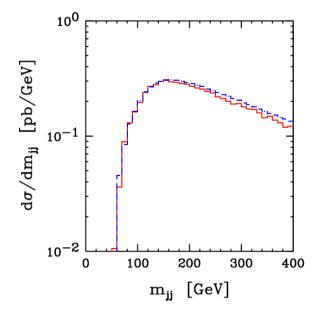
pp -> ttjj @ LHC

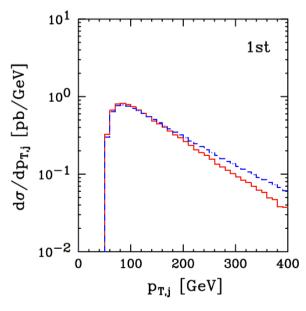
- Differential cross section
- $\rightarrow m_{\rm ff}$ size of the corrections for low p_T, shapes change

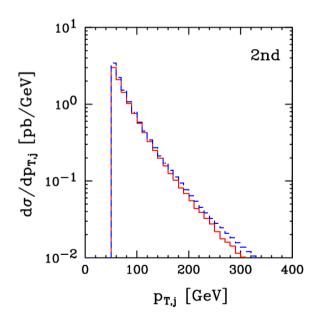
transmitted to distributions for hight p_T



 $> p_T of 1^{st} hardest & 2^{nd} hardest jet (ordered in p_T)$ altered shapes up to 39% & 28% in tails







Bevilacqua, Czakon, Papadopoulos, Worek '10

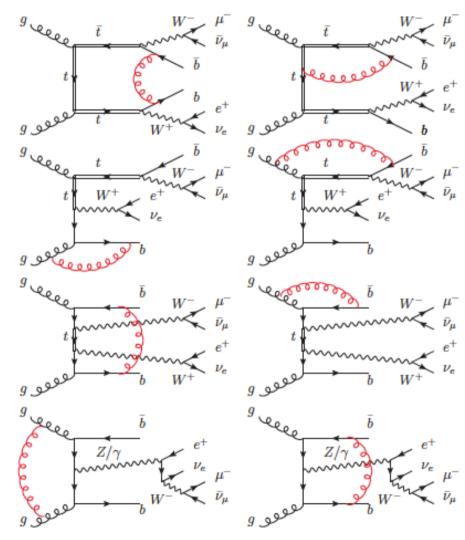
Applications

 $\gamma pp (p\overline{p}) \rightarrow WWbb$

- ☐ Complete off-shell effects @ NLO
- Double-, single- and non-resonant top contributions of the order $O(\alpha_s^3 \alpha^4)$
- ☐ Complex-mass scheme for unstable top
- ☐ W gauge bosons are treated off-shell

$$pp(p\bar{p}) \to e^+ \nu_e \mu^- \nu_\mu b\bar{b} + X$$

- Sum over helicities and color via MC
- LO + V obtained by reweighting of tree level unweighted events
- ☐ Dipole channels for subtracted real part
- ☐ Check of Ward identity for virtual part
- ☐ Cancellation of divergences
- \square α_{max} independence test for real part



- ☐ Integrated cross sections for inclusive cuts
 - $p_T(b) > 20$, GeV, $p_T(l) > 20$ GeV, $p_T(miss) > 30$ GeV
 - $|y(b)| < 4.5, |y(1)| < 2.5, \Delta R(jj) > 0.4, \Delta R(jl) > 0.4$

TeVatron	$\sigma_{ m LO}~[{ m fb}]$	$\sigma_{ m NLO}^{\alpha_{ m max}=1}$ [fb]	$\sigma_{ m NLO}^{lpha_{ m max}=0.01}$ [fb]
$anti$ - k_T	34.922 ± 0.014	35.705 ± 0.047	35.697 ± 0.049
k_T	34.922 ± 0.014	35.727 ± 0.047	35.723 ± 0.049
C/A	34.922 ± 0.014	35.724 ± 0.047	35.746 ± 0.050

LHC	$\sigma_{ m LO} \ [{ m fb}]$	$\sigma_{\rm NLO}^{\alpha_{\rm max}=1}$ [fb]	$\sigma_{ m NLO}^{lpha_{ m max}=0.01}$ [fb]	
$anti-k_T$	550.54 ± 0.18	808.46 ± 0.98	808.29 ± 1.04	
k_T	550.54 ± 0.18	808.67 ± 0.97	808.86 ± 1.03	
C/A	550.54 ± 0.18	808.74 ± 0.97	808.28 ± 1.03	



 α_{max} independence test



☐ LO & NLO scale dependence

TeVatron

$$\sigma_{LO} = 34.922^{+40\%}_{-26\%} fb$$

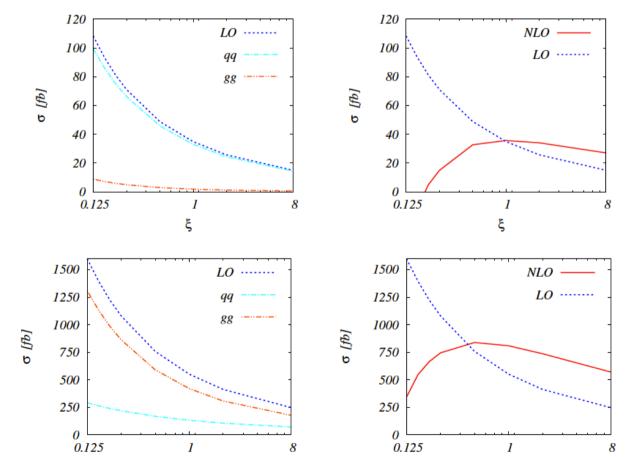
$$\sigma_{NLO} = 35.727^{-4\%}_{-8\%} fb$$

$$K = 1.023$$

LHC

$$\sigma_{LO} = 550.538 {}^{+37\%}_{-25\%} fb$$
 $\sigma_{NLO} = 808.665 {}^{+4\%}_{-9\%} fb$

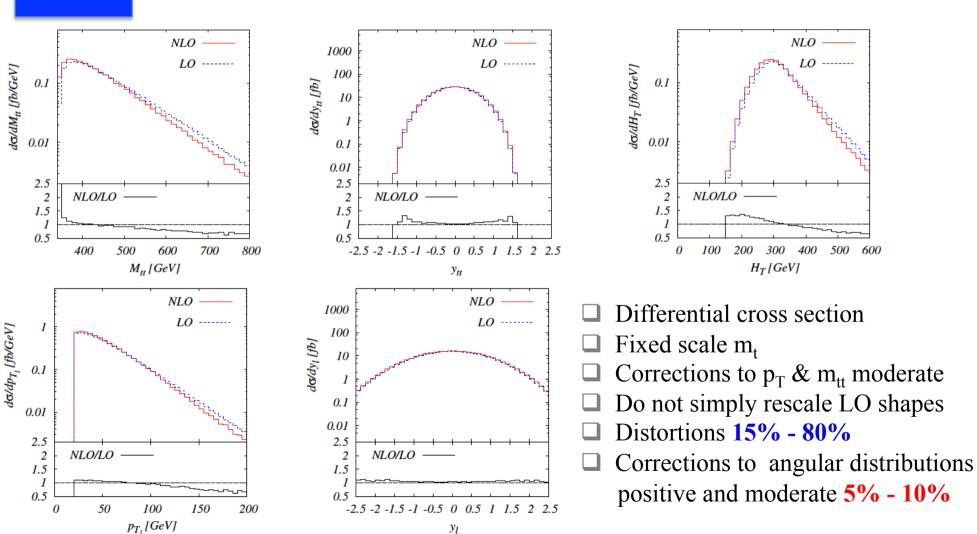
$$K = 1.47$$



Bevilacqua, Czakon, van Hameren, Papadopoulos, Worek '11

$\mathcal{P}\overline{\mathcal{P}} \rightarrow \mathcal{WWbb}$

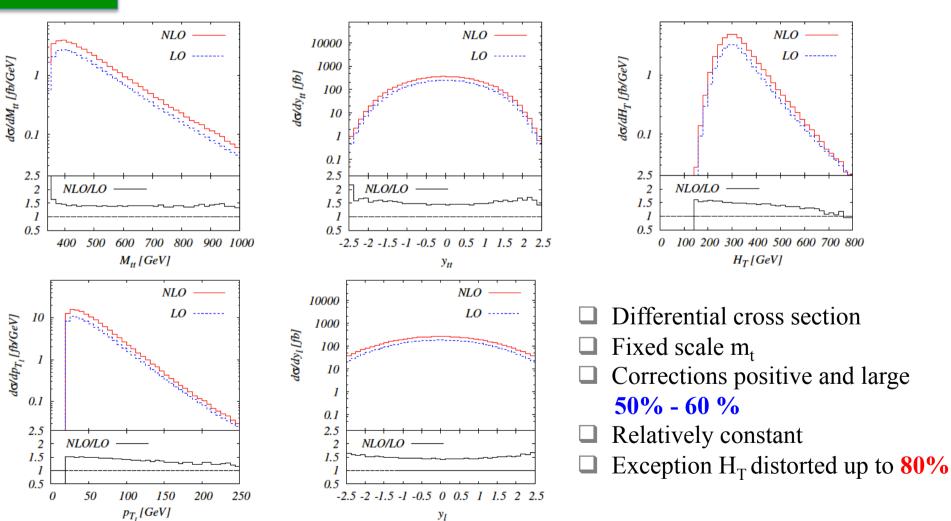
TeVatron



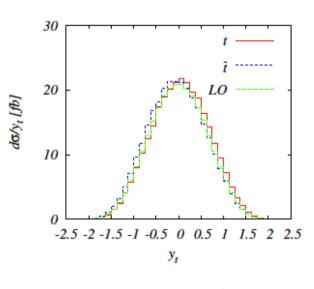
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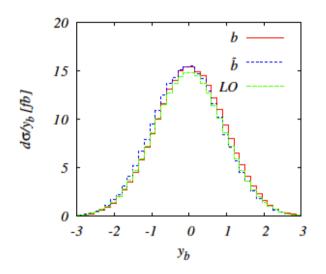
7pp -> WW66

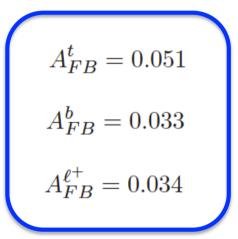


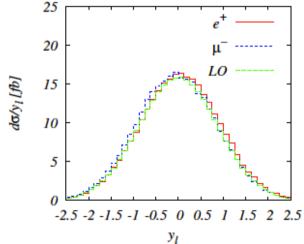


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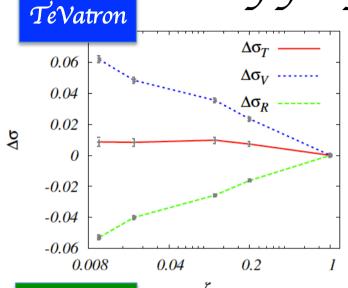






☐ Asymmetry @ TeVatron

$$A_{FB}^{t} = \frac{\int_{y>0} N_{t}(y) - \int_{y<0} N_{t}(y)}{\int_{y>0} N_{t}(y) + \int_{y<0} N_{t}(y)}$$
$$A_{FB}^{\bar{t}} = -A_{FB}^{t}$$



0.04

ζ

 $\Delta \sigma_T$

0.2

 $\Delta \sigma_V$

 $\Delta \sigma_R$ -----

LHC

0.06

0.04

0.02

-0.02

-0.04

-0.06

0.008

0

ΔG

- ☐ Size of the non-factorizable corrections
- Full result versus narrow width approximation
- Rescaling coupling tWb by some large factors
- +1.0% TeVatron and -1.2% LHC for inclusive cuts
- $lue{}$ Dependence of NLO cross section and individual contributions on rescaling parameter ζ

$$\Gamma_{rescaled} = \zeta \Gamma_t$$

$$\Delta \sigma_i(\zeta) = (\sigma_i(\zeta) - \sigma_i(\zeta = 1))/\sigma_T(\zeta = 1)$$

Bevilacqua, Czakon, van Hameren, Papadopoulos, Worek '11

Summary & Outlook

- □ ttbb, WWbb completed by two groups, Permille level agreement in both cases!
- □ ttjj

See Stefan Kallweit talk on Saturday

■ HELAC-NLO

- ➤ Complete tool at NLO built around **HELAC-PHEGAS**: **HELAC-1LOOP**, **CUTTOOLS**, **ONELOOP** & **HELAC-DIPOLES**
- Much wider study for ttjj: variation of the center of mass energy, jet algorithms, cone size in jet algorithm, transverse momentum cuts, ...
- > Other processes from NLO Wishlist under attack
- Constant improvements in speed and functionality

http://helac-phegas.web.cern.ch/helac-phegas/